CONTENTS

Regular Papers

217 62.5-Gbit Optical Receiver Using a CMOS-Compatible Si Avalanche Photodetector
   | Hyo-Seon Kang, Myung-Jae Lee, and Woo-Young Choi

221 Hydrogen Sensor Based on a Palladium-Coated Long-Period Fiber Grating Pair
   | Young-Ho Kim, Myoung-Jae Kim, Min-Su Park, Jae-Hyung Kang, Byeong-Ha Lee, and Kwang-Taek Kim

225 Optimization of 40 Gbps WDM Systems Using Super-Gaussian FZP Pulses
   | Jong-Hyung Lee, Deu-Hyun Han, Yong-Jae Lee, and Byeong-Yoon Choi

232 PDM Tolerance of 10 Gbps Modulated Signals due to SOA-Induced Grit in SOA Booster Amplifiers
   | Hodook Jung, Kyungseok Kim, Jaeheon Lee, and Jechi Jeong

240 Derivation of the Poschini and Shepp's Joint-Characteristic Function for the First-and Second-Order Polarization-Dependent Dispersion Using the Fokker-Planck Method
   | Jae-Sung Lee

244 Performance Evaluation of a Tunable Dispersion Compensator based on Strain-Chirped Fiber Bragg Grating in a 40 Gbps Transmission Link
   | Chul Han Kim, Junseok Bae, Kwansik Lee, and Seng Bae Lee

249 An Electronic Domain Chromatic Dispersion Monitoring Scheme Insensitive to OSNR Using Kurtosis
   | Kyung-woo Kim, Jaeheon Lee, Wonwoo Chung, and Sungchul Kim

255 Side-Coupled Asymmetric Plastic Optical Fiber Coupler for Optical Sensor Systems
   | Kwang-Taek Kim, Doo-Hyung Lee, Young-Kwan Hyun, Ki-Bum Hong, Kizgon Kim, So-Jung Bak, Deukyung Kim, and Hwan-joon Choi

262 Optical Design of a Compact Imaging Spectrometer for STSAT3
   | Jan Ho Lee, Tae Seong Kang, Ho-Soon Yang, and Seung-Woo Rhee

269 Hologram Generation of 3D Objects Using Multiple Orthographic View Images
   | Min-Su Kim, Gunho Bae, Sungnam, Nam Kim, and Jae-Hyeong Park

275 Elemental Image Synthesis for Integral Imaging Using Phase-shifting Digital Holography
   | Min-Seok Jeong, Nam Kim, and Jae-Hyeong Park

281 3-D Surface Profile Measurement Using an Acousto-optic Tunable Filter Based Spectral Phase Shifting Technique
   | Dae-je Kim and Yong-Jae Cho

288 Analysis of Temperature Effects on Reaman Silicon Photonic Devices
   | Wonsuk Kim and Dongwook Park

296 270-W 15-kHz MOPA System Based on Side-pumped Rod-type Nd:YAG Gain Modules
   | Yong-Ho Choo, Myoung-Yul Yang, Kwang-Hoon Ko, Gwon Lim, Jae-Min Han, Hyun-Min Park, Tai-Kyu Kim, Sunho Rok, and Do-Yeong Jeong

   | Hyuk-Kyu Song and Ming C. Wu

309 Comparison of Sapphire and Germanium Fibers for Erbium-Yag Lithotripsy
   | Joo Lee, Jo-Wook Yoo, Young-Dae Kim, Jeong Hyun Kim, Robert T. Ryan, Poo M.H. Tchikman, and A.J. Welck

314 Complete Tunneling of Light via Local Barrier Modes in a Composite Barrier with Metamaterials
   | Kwang-Youn Kim and Sunhwa Kim

319 Phase Contrast Projection Display Using Photopolymer
   | Mei-Lun Pan, Nam Eun, and Jae-Hyeong Park

Special Issue: WDM-PON
The special issue papers were recommended by the Editorial Board of Photonics Division to promote fiber-to-the-home including WDM-PON technology, and selected from all the papers presented in the Photonics Conference 2003.
The papers were accepted for publication in the Journal of the Optical Society of Korea after standard review processes.

325 WDM-PON Based on Wavelength Locked Fabry-Perot LDs
   | Chang-Hye Lee and Sil-Cu Mun

337 RSOA-Based Wavelength-Reuse Gigabit WDM-PON
   | Byung-Whi Kim

346 Performance Analysis of Metro WDM WDM Star Network based on An Arrayed Waveguide Grating
   | Syed Abdul Masum Khan, Seung-Sup Shim, Hyo-Sik Yang, and Noemin Zafar Azemi

351 Ultra-Dense WDM PON with 12.5-GHz Spaced 256 Channels
   | Jae-Nam Yim, Gyu-San Hwang, Jae-Seung Lee, Kyung-Hoon Seo, Hyun-Jae Lee, and Jeeseo Ko

355 Fabrication of 2.5-Gbs Burst-mode Receiver and Its Full Compliance to GPON
   | Min-Soo Lee, Byung-Tak Lee, Yong-Doo Kim, and Dong-Soo Lee

359 Bidirectional 12.5-Gbs WDM-PON with Broadcasting Function Using a Fabry-Perot Laser Diode and RSOA
   | Tae-Kyung Piao, Hyung-Seung Kim, Yong-Yuk Won, and Sang-Koon Han

364 Erratum

365 Information for Contributors

366 Optical Society of Korea, Copyright Transfer Agreement
Elemental Image Synthesis for Integral Imaging Using Phase-shifting Digital Holography

Min-Ok Jeong, Nam Kim, and Jae-Hyeung Park*

School of Electrical & Computer Engineering, Chungbuk National University
410 SungBong-Ro, Heungduk-Gu, Cheongju-Si, Chungbuk, 361-763, Korea

(Received October 8, 2008 : revised December 3, 2008 : accepted December 3, 2008)

We propose a method generating elemental images for the integral imaging using 4-step phase-shifting digital holography. Phase shifting digital holography is a way recording the digital hologram by changing the phase of the reference beam and extracting the complex field of the object beam. Since all 3D information is captured by phase-shifting digital holography, the elemental images for any specifications of the lens array can be generated from single phase-shifting digital holography. In experiment, phase-shifting is achieved by rotating half- and quarter- wave plates and the resultant interference patterns are captured by a 3272×2469 pixel CCD camera with 27 μm×27 μm pixel size.

Keywords : Digital holography, Phase shifting holography

OCIS codes : (090.0090) Holography; (100.0100) Image processing; (090.2880) Holographic interferometry; (130.0130) Integrated optics

I. INTRODUCTION

Imaging of three-dimensional(3D) objects is important for various applications including display. Holography is one of the useful 3D imaging methods since it can record full 3D information of the object. Photographic recording and optical reconstruction requiring chemical processing, however, have been problematic for conventional holography [1]. Overcoming the disadvantages of the conventional holography, digital holography has been proposed recently. The digital holography uses Charge Coupled Devices(CCD) for recording holographic patterns and reconstructs 3D object images by numerical Fresnel propagation. The whole process is simple because of the use of the CCD instead of the holographic film, the processing time is within the practical range due to the development of the computing power. With these advantages, the digital holography has been used in various practical applications including spatial filtering, inspection of optical systems, tests of laser beams and reconstruction of 3D images, etc [2-5].

Integral imaging is one of the attractive auto-stereoscopic 3D display techniques. Using a lens array, integral imaging can display full parallax, auto-stereoscopic 3D images continuously within a viewing angle without any eye fatigue [6-9]. One problem of integral imaging is that the lens array should be the same for the pickup and the display process. If the lens array of different specifications is used for the pickup and display process, the 3D image will be distorted severely. Generally, it is hard to expect that all the integral imaging display systems have lens arrays of the same specifications since it should vary according to the display environment and desired applications. Many variations of the integral imaging display for enhancing the viewing angle, available depth range and the viewing resolution also require special lens arrays or elemental images matched to their own purpose [10-14]. It is obviously not practical to repeat capturing the elemental images for the different types of display systems. Therefore it is necessary to devise a method to capture the 3D information of the object by a single pickup process and process the captured 3D information to synthesize the specific elemental images for various different display systems. Although a scaling technique of the elemental images has been reported to deal with this problem, it does not provide full freedom in the lens array specification [7-9].
Hologram is a good alternative for direct capturing. Since a hologram captures full 3D information about the object, it is possible to synthesize the desired elemental images from the single hologram of the object. In this report, we propose a method generating elemental images for integral imaging using phase-shifting digital holography. The proposed method obtains the complex field of the object wave at the CCD plane using 4-step phase shifting digital holography. By propagating the obtained complex field numerically assuming a virtual lens array, the proposed method synthesizes the elemental images. Since we can set the specifications and the location of the lens array arbitrarily in the numerical imaging, we can generate the elemental images corresponding to any integral imaging display systems from single capture of phase shifting digital holography.

The proposed method is closely related to Javidi et al.’s method [15]. Javidi et al. also synthesizes the elemental images from the object complex field obtained by the 4-step phase shifting digital holography. Their method, however, generates the elemental image by projecting and accumulating the object plane images reconstructed at various depths which sweep the object space onto the elemental image plane assuming the pinhole array based on ray optics. Therefore it is not possible to synthesize the elemental images for the lens array of the arbitrary specifications including the aperture size of the elemental lens, and irregular curvature of the elemental lens surface. The proposed method can deal with the lens array of any specifications since it is fully based on wave optics and does not use a pinhole approximation. Kim et al [16], also proposed an elemental image generation method from the hologram. The elemental image plane, however, is fixed in their method and the elemental image synthesis for arbitrary lens array locations and specifications, which is easily realized in the proposed method, is not presented.

In the following, we explain the principle of the proposed method and provide the experimental results for verification.

II. PRINCIPLE

The first step of the proposed method is the capture of the object complex field using phase-shifting digital holography. Figure 1 is the coordinate system for the numerical reconstruction of phase-shifting digital holography. The object wave at the CCD camera plane \( U(x,y) \) is represented as a Fresnel Transform of the complex field at the object plane \( U_o(x', y') \) by

\[
U(x,y) = \int \int U_o(x', y') \exp \left[ i \frac{(x-x')^2 + (y-y')^2}{2z_o} \right] dx' dy'
\]

\[= A(x,y) \exp [i \phi(x,y)] \tag{1} \]

where \( A(x,y) \) is the amplitude of the object wave recorded by the CCD camera, and \( \phi(x,y) \) is the phase of the object wave.

The coordinate system for the numerical reconstruction is shown in Fig. 1. The object complex field \( U_o(x', y') \) is represented as a Fresnel Transform of the elemental image plane at the CCD camera. A plane reference wave whose complex amplitude is represented by \( U_r(x,y) = A_r(x,y) \exp [i \phi(x,y)] \) and \( \exp [i \phi(x,y)] \) is superimposed upon the object wave. Since the interference patterns detected by CCD camera have only the intensity value, the resultant interference intensity recorded by the CCD camera is given by

\[I(x,y; \alpha) = |U_r(x,y)+U_o(x,y)|^2\]

\[= |A_r(x,y)|^2 + |A_o(x,y)|^2 + 2A_r(x,y)A_o(x,y) \cos[\phi(x,y) - \phi(x,y) - \alpha]\]

\[= |A_r(x,y)|^2 + |A_o(x,y)|^2 + 2A_r(x,y)A_o(x,y) \cos[\phi(x,y) - \phi(x,y) - \alpha] \tag{2} \]

where \( \alpha \) is the mean phase difference between the object and the reference waves. By using the phase-shifting procedure in which the reference phase is shifted by a step of \( \pi/2 \), we can derive the object phase \( \phi(x,y) \) such that

\[\phi_r(x,y) = \phi(x,y) - \phi(x,y) = \tan^{-1} \left[ \frac{I(x,y; \pi/2) - I(x,y; 3\pi/2)}{I(x,y; 0) - I(x,y; \pi)} \right] \tag{3} \]

The object amplitude, \( A_o(x,y) \), can be calculated from the following equation,

\[A_o(x,y) = A_r(x,y) \exp [i \phi(x,y)] \tag{4} \]

where the argument of the cosine function in the denominator is directly obtained from Eq. (3) [17].

If we use a plane reference wave, the amplitude \( A_r(x,y) \) and the phase \( \phi_r(x,y) \) of the reference wave can be considered as constants. Hence, by using Eqs. (3) and (4), we can derive the complex amplitude of the object wave \( U \) at the CCD plane [18]. Since the complex object field is obtained at the CCD plane, we can calculate the object field at any plane by numerically propagating the obtained object field.

\[U(x,y) = \int \int U_o(x', y') \exp \left[ i \frac{(x-x')^2 + (y-y')^2}{2z_o} \right] dx' dy' \]

\[= A(x,y) \exp [i \phi(x,y)] \tag{1} \]
for simplicity, but its extension to the lens array is straightforward. In order to obtain the elemental image which is denoted by $U_d$ in Fig. 3, we first calculate the optical field $U_i$ at the front side of the virtual lens by numerically propagating $U$ by a distance $z_i$ using Eq. (5). The optical field $U'_i$ right behind the lens is given by

$$U'_i(X,Y)=U_i(X,Y,z_i)P(X,Y)\exp\left[-\frac{ik(x-x_i)^2+(Y-y_i)^2}{2f}\right]$$

where, $P(X,Y)$, $(x_o,y_o)$, and $f$ are the pupil function, the transverse location and the focal length of the virtual lens, respectively. The elemental image is finally obtained by numerically propagating $U'_i$ over distance $-d$ [19].

$$U_e(X,Y,-d)=\int\int U'_i(X,Y)\exp\left[-\frac{ik(x-x')^2+(Y-y')^2}{2d}\right]dx\,dy$$

Note that the sign of the exponential is different from Eq. (5) because of the opposite direction against $Z$. At this time, $d$ can be determined by the lens’ law, such as

$$\frac{1}{f} = \frac{1}{z_o-z_i} + \frac{1}{d}$$

$$d = \frac{(z_o-z_i)f}{(z_o-z_i)-f}$$

where $z_o$ is the distance between the plane of CCD camera and the object and $z_i$ is the distance between the $U_i$ plane and the $U$ plane in Fig. 3. By repeating this process for all elemental lenses, the whole elemental image can be obtained.

The proposed method can generate the elemental images for any lens array specifications including the focal length, elemental lens size, number of the elemental lenses and elemental lens arrangement. The resolution and the viewing angle of the elemental image, however, are limited by the sampling interval and the viewing angle of the complex object field captured at the CCD plane. The viewing angles of the object field captured at the CCD camera plane, $\theta_i$ and $\theta_o$ are given by

$$\theta_i = 2\tan^{-1}\frac{\delta_x N_i/2}{z_i}$$

$$\theta_o = 2\tan^{-1}\frac{\delta_x N_o/2}{z_o}$$

where pixel size of the CCD camera are $\delta_x\times\delta_y$ and the number of pixels are $N_x\times N_y$.

The sampling interval varies at each plane. The sampling interval of $U_i$ plane, $\Delta x$ and $\Delta y$, are given by
FIG. 4. (a) Arrangement for recording 4-step phase shifting digital hologram: SF, spatial filter; M, mirror; P1, P2, polarizer; NDF, neutral density filter; BS1, BS2, beam splitter; RP1, RP2, RP3, retardation plate.

(b) 4-step phase shifting using retardation plates.

\[ \Delta \phi = \frac{\Delta z_i}{\delta_N x} \]
\[ \Delta \theta = \frac{\Delta z_i}{\delta_N y} \]

where the sampling intervals at the CCD plane are \( \delta_x \) and \( \delta_y \) [20]. Then the final sampling intervals of the elemental images at the \( U_d \) plane are given by

\[ \Delta \phi = \frac{\lambda d}{\Delta N x} \]
\[ \Delta \theta = \frac{\lambda d}{\Delta N y} \]

III. EXPERIMENTAL RESULTS

Figure 4(a) shows the experimental setup for recording a 4-step phase-shifting hologram based on Mach-Zehnder interferometer [21]. The Nd:YAG laser beam (532nm) passes through the NDF (Neutral Density Filter), the polarizer and the \( \lambda/2 \) wave plate and then divided by a beam splitter into object and reference beams. The reference beam passes through the other \( \lambda/2 \) wave plate and \( \lambda/4 \) wave plate. The object beam illuminates the object after reflection by two mirrors. These two beams are put together by another beam splitter and their interference pattern is captured by the CCD camera. We used a coin located at 56.5 cm from the CCD as an object. The CCD camera used in the experiment has 3272×2469 pixels with 27 \( \mu \)m×27 \( \mu \)m pixel size. The phase-shifting of the reference beam is achieved by rotating the \( \lambda/2 \) and \( \lambda/4 \) wave plates, RP2 and RP3 in Fig. 4(a). By aligning different slow and fast axes with the direction of the polarization, 4-step phase values \( \alpha = 0, \pi/2, \pi \) and \( 3\pi/2 \) can be produced with high accuracy as shown in Fig. 4(b) [17]. Figure 5 shows the interference patterns captured by the CCD for 4 phase shifts. From these 4 interference patterns, the complex object field is calculated by Eqs. (3) and (4) as shown in Fig. 6. The calculated values of the view angles of CCD camera plane are \( \theta_x = 0.017^\circ \) and \( \theta_y = 0.013^\circ \) respectively by Eq. (10).

Using the calculated complex object field, the elemental images for various parameters are synthesized. Figure 7 shows an elemental image generated assuming only one elemental lens. The parameters used in generating Fig. 7 are \( f = 4 \) mm, \( z_i = 6 \) cm, \((x_o, y_o) = (2 \text{ mm}, 0)\), \( W_x = W_y = 2 \) mm, \( d = 4.032 \) mm where \( W_x \) and \( W_y \) are the lens size. We can see that a clear image of the object is reconstructed by the proposed method. The object image size measured in Fig. 7 is 597×666 pixels or 0.10834 mm×0.12087 mm considering the sampling interval calculated by Eq. (12). Note that because the actual size of the coin is about 1.5 cm×1.5 cm if we get the elemental...
FIG. 7. The reconstructed elemental image in \((x_o, y_o) = (0, 0)\), \(W_x = W_y = 2 \text{ mm}\), \(f = 4 \text{ mm}\), \(z_l = 6 \text{ cm}\); \(d = 4.0329 \text{ mm}\), full image size = 0.59381 mm \(\times\) 0.44808 mm.

FIG. 8. The reconstructed elemental images in \(z_l = 7 \text{ cm}\), \(f = 5 \text{ mm}\), \(W_x = W_y = 3 \text{ mm}\); \(d = 5.05263 \text{ mm}\) (a) \((x_o, y_o) = (-3 \text{ mm}, 3 \text{ mm})\), (b) \((x_o, y_o) = (0, 0)\) and (c) \((x_o, y_o) = (3 \text{ mm}, -3 \text{ mm})\).

image optically with actual elemental lens of the same parameters, then its size will be \[\text{actual object size} \times d / (z_o - z_l) = [0.015 \text{ m}] \times 0.0040329 \text{ m} / (0.565 \text{ m} - 0.06 \text{ m}) = 0.1198 \text{ mm},\]
which is in a good match with the measured value in Fig. 7. Therefore, we can confirm that the size of the elemental image produced by the proposed method is realistic for given parameters. The disparity between the elemental images is also verified experimentally. Figure 8 shows the generated elemental images for 3 different elemental lens locations \((x_o, y_o)\), i.e. \((0 \text{ mm}, 0 \text{ mm})\), \((-3 \text{ mm}, 3 \text{ mm})\), and \((3 \text{ mm}, -3 \text{ mm})\). The dotted line circle shown in Fig. 8 is the position of the elemental image when the elemental lens is at \((0 \text{ mm}, 0 \text{ mm})\). Since the separation between the elemental lenses is \(3 \sqrt{2} \text{ mm}\), the disparity should be \(\text{lens' separation} \times d / (z_o - z_l) = [3 \text{ mm}] \times 5.05263 \text{ mm} / 495 \text{ mm} = 0.0433 \text{ mm}\). The disparity measured experimentally from Figs. 8(a) and (b) is \((160 \text{ pixel}, 148 \text{ pixel})\), or \(\sqrt{(160^2 + 148^2) \times 10^{-4} \text{ mm} = 0.0425 \text{ mm}}\), considering the sampling interval. This measured disparity value 0.0423 mm is well matched to the expected value 0.0433 mm, confirming the proposed method produces the elemental images of correct disparity. Figures 9-10 shows the generated elemental images for various elemental lens focal lengths \(f\), and longitudinal locations \(z_l\). The sizes of all images in Fig. 9 and Fig. 10 are normalized, considering their sampling interval, for easier comparison. It is observed from Figs. 9-10 that the elemental images are generated correctly with given parameters by the proposed method. Finally, Fig. 11 shows the generated full set of elemental images assuming a lens array of the parameters \(f = 4 \text{ mm}\), \(z_l = 5 \text{ cm}\), \(W_x = W_y = 2 \text{ mm}\). We can see the elemental images are generated clearly with the proper disparity and realistic sizes, confirming the feasibility of the proposed method. Although only \(5 \times 3\) elemental images are shown in Fig. 11, it is straightforward to increase the number of the elemental images. The circular noise seen in each elemental image might be produced by the unwanted reflection from the optical components used in the experimental setup. It is believed that this noise will be removed by a more refined experiment.

IV. DISCUSSION AND CONCLUSION

In this paper, we proposed a method generating elemental images for integral imaging using the 4-step
phase shifting digital hologram. After the complex field of the object is obtained by the 4-step phase shifting digital hologram, we propagate the complex object field numerically assuming a virtual lens array in the reconstruction space, resulting in the elemental images. We can control the parameters about elemental lens’ focal length $f$, propagated distance of obtained hologram $z_l$, elemental lens’ transverse location $(x_0, y_0)$ and elemental lens’ horizontal and vertical length $W_x$ and $W_y$. Since the hologram contains full 3D information of the object, we can generate the elemental images of the object for arbitrary lens array. Therefore from single acquisition of the object hologram, elemental images for any integral imaging display system can be synthesized regardless of its specifications. The result verifies the principle of the proposed method.

ACKNOWLEDGMENT

"This work was supported by the grant of the Korean Ministry of Education, Science and Technology" (The Regional Core Research Program / Chungbuk BIT Research-Oriented University Consortium)

REFERENCES